

Brownian motion with limited local time

M.Savov (joint work with M.Kolb)

In this talk we discuss one dimensional Brownian motion with zero drift $(B_s)_{s \geq 0}$ and local time at zero defined as

$$L_t = \lim_{\epsilon \rightarrow 0} \frac{1}{2\epsilon} \int_0^t \mathbf{1}_{\{|B_s| \leq \epsilon\}} ds.$$

We know that L_t is a non-decreasing stochastic process which increases only on $\{s \geq 0 : B_s = 0\}$.

We are interested in models in which a deterministic, continuous and increasing function $f : [0, \infty) \mapsto (0, \infty)$ restricts the rate of return to zero L_t , i.e. we shall investigate the conditional measures

$$\mathbb{P}_t(\cdot) = \mathbb{P}(\cdot | L_s \leq f(s), 0 \leq s \leq t).$$

We show for a class of functions f that the Brownian motion with restricted local time exists and describe its properties, i.e. show that

$$\lim_{t \rightarrow \infty} \mathbb{P}_t(\cdot) = \mathbb{Q}(\cdot)$$

and discuss \mathbb{Q} .

Example: When $f(t) \equiv 0$ then we have that \mathbb{Q} is the law of the three dimensional Bessel process started from 0. This corresponds to the extreme case when the Brownian motion is not allowed to return to zero.

Previous work: This problem is introduced and studied by I. Benjamini and N. Berestycki. Their contributions are the following:

1 Always $(\mathbb{P}_t)_{t \geq 0}$ is a tight sequence of measures

2 If

$$I(f) := \int_0^\infty f(s) \frac{ds}{s^3} < \infty$$

then every **possible** limit \mathbb{Q} of \mathbb{P}_t is the law of a transient process namely

$$\mathbb{Q} \left(\lim_{t \rightarrow \infty} |X_t| = \infty \right) = 1.$$

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They proposed the following open problems:

- 1 If $I(f) < \infty$ then is it true that $\mathbb{P}_t \rightarrow \mathbb{Q}$ and can we explicitly describe \mathbb{Q} ?
- 2 If $I(f) = \infty$ then is it true that $\mathbb{P}_t \rightarrow \mathbb{Q}$? Is the limiting process recurrent to zero ?
- 3 If $I(f) = \infty$ can we describe the **decreasing to zero** functions $w(t)$ such that

$$\lim_{t \rightarrow \infty} \mathbb{Q}(L_t \leq w(t)f(t)) = 1.$$

$\mathcal{D} = \{w : \lim_{t \rightarrow \infty} \mathbb{Q}(L_t \leq w(t)f(t)) = 1\}$ is called the **entropic repulsion** envelop. *Entropic repulsion* is the phenomenon whence the easiest way for a process to satisfy an imposed condition is to satisfy even more stringent condition.

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The starting point of our study is the following immediate observation

$$\mathcal{O}_t = \{\tau_s \geq g(s), s \leq t\} = \{L_s \leq f(s), s \leq g(t)\},$$

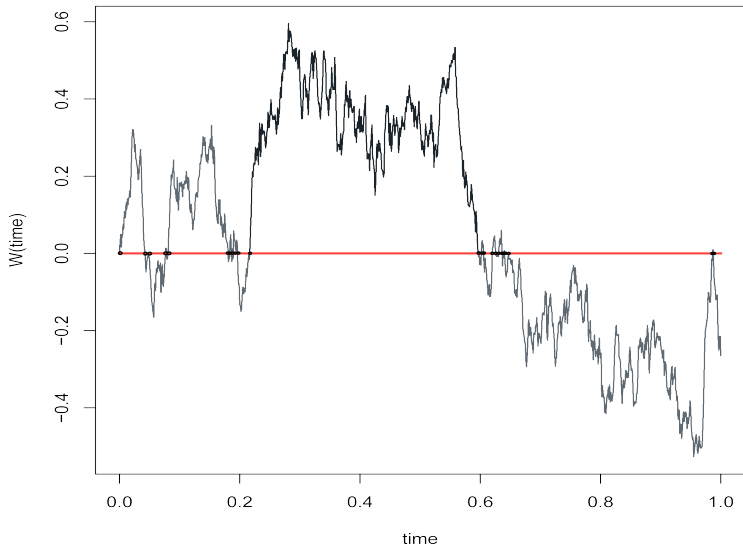
where $\tau = L^{-1}$ is the right inverse of L and $g = f^{-1}$.

We know that $(\tau_t)_{t \geq 0}$ is a stable subordinator (increasing Lévy process) with index $1/2$, namely

$$\mathbb{E} \left[e^{-\lambda \tau_t} \right] = e^{-t\lambda^{\frac{1}{2}}}, \lambda > 0.$$

Also $\tau_t = \sum_{s \leq t} \Delta \tau_s$. The jumps $\Delta \tau$ are the length of the excursions away from zero.

A Typical Path of Brownian Motion



- 1 It suffices to study the limit

$$\lim_{t \rightarrow \infty} \mathbb{P}(\tau \in \cdot | \mathcal{O}_t) = \mathbb{Q}(\tau \in \cdot)$$

and if the limit process τ exists we then conditionally on a path realization of τ we splice Brownian excursions between the end points of each jump of τ .

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Some ideas for the proof:

- 1 Using the Markov property we have that

$$\begin{aligned}\mathbb{P}(\tau_a \in dx | \mathcal{O}_t) &= \frac{\mathbb{P}(\tau_a \in dx, \mathcal{O}_t)}{\mathbb{P}(\mathcal{O}_t)} = \\ \frac{\mathbb{P}(\tau_a \in dx, \mathcal{O}_a \cap \mathcal{O}_t)}{\mathbb{P}(\mathcal{O}_t)} &= \frac{\mathbb{P}(\mathcal{O}_t | \tau_a = x, \mathcal{O}_a) \mathbb{P}(\tau_a \in dx, \mathcal{O}_a)}{\mathbb{P}(\mathcal{O}_t)} \\ \frac{\mathbb{P}(\mathcal{O}_{t-a}^{x,a})}{\mathbb{P}(\mathcal{O}_t)} \mathbb{P}(\tau_a \in dx, \mathcal{O}_a) &.\end{aligned}$$

- 2 Above $g^{x,a}(s) = g(s+a) - x$ and

$$\mathcal{O}_t^{x,a} = \{\tau_s > g^{x,a}(s), s \leq t\}.$$

- 3 However, to prove recurrence, we need *uniform* asymptotic of $\mathbb{P}(\mathcal{O}_{t-a}^{x,a})$, so as to apply DCT in

$$1 = \mathbb{P}(\tau_a > g(a) | \mathcal{O}_t) = \int_{g(a)}^{\infty} \frac{\mathbb{P}(\mathcal{O}_{t-a}^{x,a})}{\mathbb{P}(\mathcal{O}_t)} \mathbb{P}(\tau_a \in dx, \mathcal{O}_a)$$

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Intuition behind the asymptotic of \mathcal{O}_t :

- Since τ_s is non-decreasing a jump Δ exceeding $g(t)$ at time $T < t$ will ensure that $\mathcal{O}_t = \{\tau_v > g(v), v \leq t\} = \mathcal{O}'_T \cap \{T \leq t\}$.
- $\mathcal{O}'_s = \{\tau'_v > g(v), v \leq s\}$ and τ' is obtained from τ by truncating all jumps larger than $g(t)$.
- Then the total probability formula gives

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- Therefore

$$\phi'(t) = \frac{2K}{\sqrt{g(t)}} \phi(t) + R_g(t) \phi(t)$$

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$$\Phi(t) = \Phi(t_0)e^{\int_{t_0}^t \frac{2K}{\sqrt{g(s)}} ds + \int_{t_0}^t R_g(s) ds}$$

- When $\int_{t_0}^{\infty} R_g(s) ds < \infty$ the asymptotic is entirely determined by $\int_{t_0}^t \frac{2K}{\sqrt{g(s)}} ds$. When τ is stable with index $\alpha = 1/2$ this is the case when $\liminf_{t \rightarrow \infty} g(t)/(t^2 \ln^{\frac{8}{5} + \epsilon}(t)) = \infty$.
- When $\int_{t_0}^{\infty} \frac{2K}{\sqrt{g(s)}} ds < \infty$ we have our criterion $I(f) < \infty$ and $\Phi(\infty) = \int_0^{\infty} \mathbb{P}(\mathcal{O}_s) ds < \infty$.
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If $\int_{t_0}^{\infty} R_{g_i}(s) ds < \infty, i = 1, 2$. Then the asymptotic of $\mathbb{P}(\tau_s > g_i(s), s \leq t)$ are directly comparable and tightly related if furthermore

$$\int_{t_0}^{\infty} \left| \frac{1}{\sqrt{g_1(s)}} - \frac{1}{\sqrt{g_2(s)}} \right| ds < \infty.$$

This is how we prove **uniformity** in the asymptotics of $\mathbb{P}(\tau_s > g(s+a) - x, s \leq t) = \mathbb{P}(\tau_s > g^{x,a}(s), s \leq t)$

Under

$$I(f) = \int_1^\infty \frac{f(s)}{s^{\frac{3}{2}}} ds < \infty$$

we have that the limiting measure corresponds to the process:

- 1 Sample from the random variable X with density $\phi^{-1}(\infty)\mathbb{P}(\mathcal{O}_s) ds$
- 2 Conditional on $X = s$ run a Brownian motion with the restriction $\{L(v) \leq f(v), v \leq s\}$.
- 3 At the moment s choose with equal probability the value $Y = \pm 1$.
- 4 We choose independent Bessel process $B^{(3)}$ and from time s we attach $YB^{(3)}$ (the attachment of this process plays the role of an excursion away of zero of infinite length at time X).

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- 2 Under \mathbb{Q} the process is recurrent at zero. *In this case $\Phi(\infty) = \infty$ and infinite excursion away from zero is not attached*
- 3 For a function $w(t) \downarrow 0$ we have that

$$\lim_{t \rightarrow \infty} \mathbb{Q}(L(t) \leq f(t)w(t)) = 0 \iff \lim_{t \rightarrow \infty} \int_t^{f\left(\frac{g(t)}{w(t)}\right)} \frac{1}{\sqrt{g(s)}} ds = 0.$$

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$$\lim_{t \rightarrow \infty} \mathbb{Q}(L(t) \leq f(t)w(t)) = 0 \iff \lim_{t \rightarrow \infty} \int_t^{f\left(\frac{g(t)}{w(t)}\right)} \frac{1}{\sqrt{g(s)}} ds = 0.$$

When $I(f) = \infty$ and with $g = f^{-1}$. Then when $\liminf_{t \rightarrow \infty} g(t)/t^2 \ln^{\frac{8}{5} + \epsilon}(t) = \infty$ we have that:

- 1 $\mathbb{P}_t \rightarrow \mathbb{Q}$
- 2 Under \mathbb{Q} the process is recurrent at zero. *In this case $\Phi(\infty) = \infty$ and infinite excursion away from zero is not attached*
- 3 For a function $w(t) \downarrow 0$ we have that

$$\lim_{t \rightarrow \infty} \mathbb{Q}(L(t) \leq f(t)w(t)) = 0 \iff \lim_{t \rightarrow \infty} \int_t^{f\left(\frac{g(t)}{w(t)}\right)} \frac{1}{\sqrt{g(s)}} ds = 0.$$

Thank you!